

Nuclear Magnetic Resonance Spectroscopy

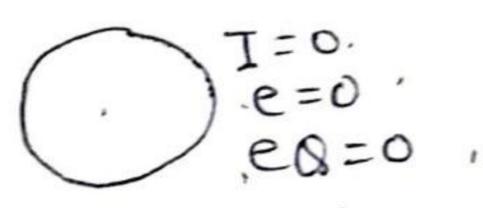
The properties necessary for understanding of NMR spectroscopy are - )

- (i) Net spin of a nucleus ( $I$ )
- (ii) Distribution of +ve charge

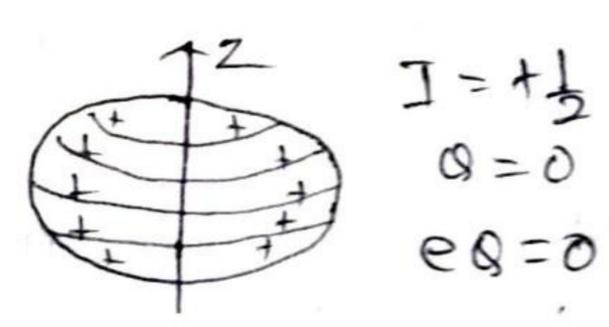
If nuclear spin ( $I$ ) is zero, there is no net spin of the nuclear charge and also its  $eQ = 0$ , ~~where~~

where,  $e$  = unit of electrostatic charge, and  
 $Q$  = measure of deviation of charge distribution from spherical symmetry.

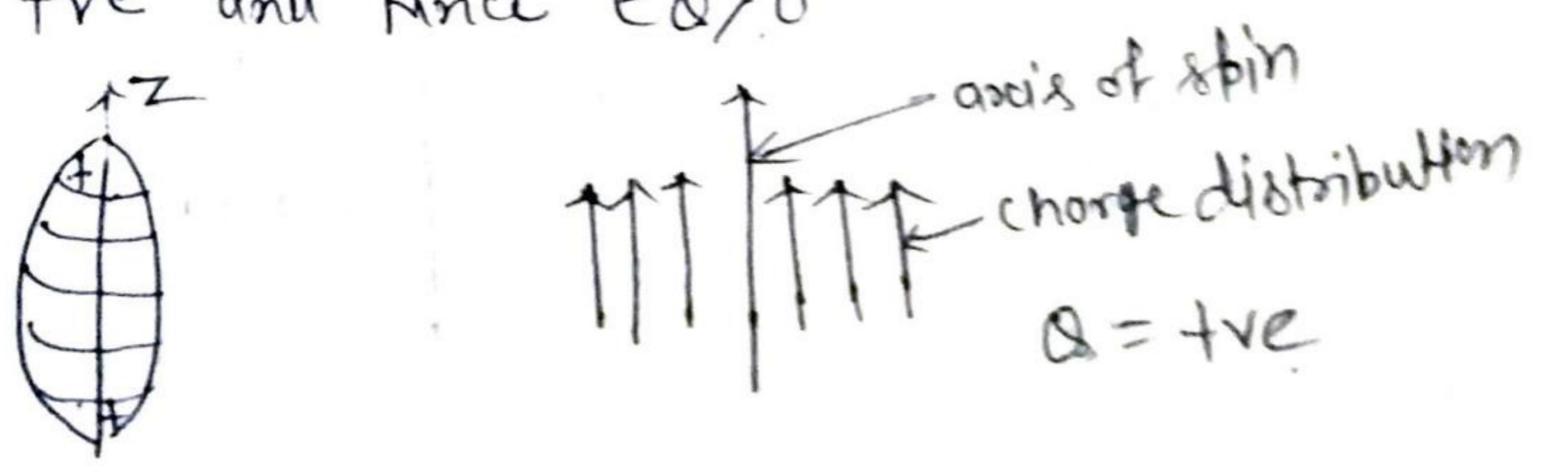
(i) For  $I = 0$ ,  $e$  is zero and hence  $eQ = 0$



(ii) When  $I = \frac{1}{2}$ , there is net spin in the nucleus, as the +ve charge distribution is spherical, so  $Q = 0$  and hence  $eQ = 0$



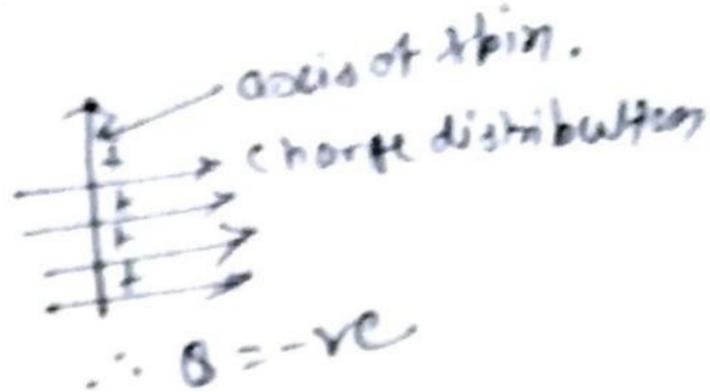
(iii) If  $I \geq 1$ , there is again net spin is zero of the nuclear charge but +ve charge distribution is non-spherical, if +ve charge distribution is elongated along the axis of spin  $Q$  is +ve and hence  $eQ > 0$



(24) If the free charge distribution is flattened about the prime axis i.e.,  $I$  to the prime axis  $(z)$  is -ve and hence  $eQ < 0$

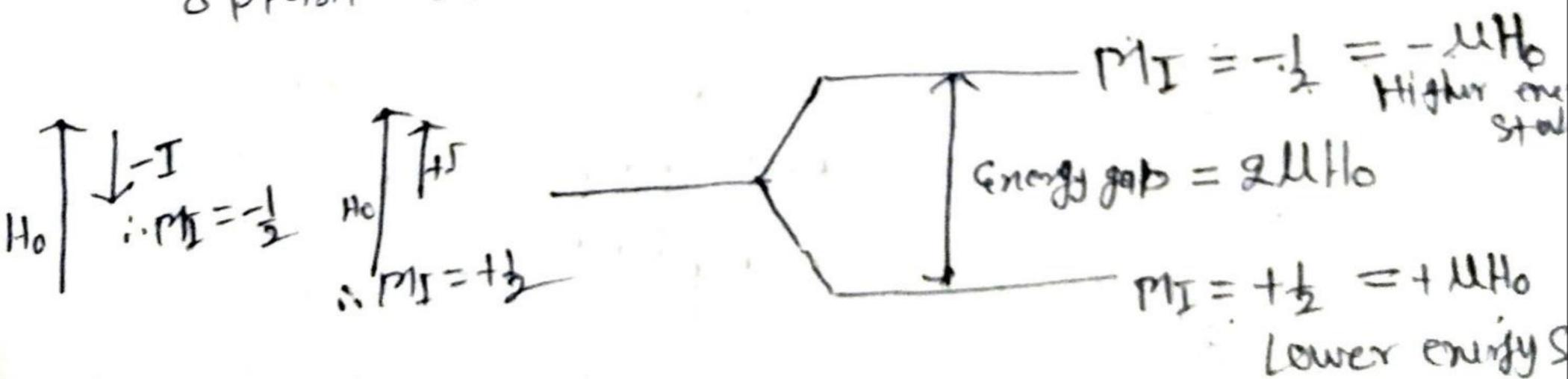


$Q = -ve$   
 $eQ < 0$



If  $I = 0$ , there is no variation of NMR spectra of such nuclei. If  $I = 1$  or  $I > 1$ , spectra is somewhat complex and only special cases are considered. So most of the NMR spectra concerned with the nuclei having  $I = \frac{1}{2}$  like  $^1H, ^{13}C, ^{19}F, ^{31}P$  etc.

When  $I = \frac{1}{2}$ , there can be two allowed orientations of the nuclear magnetic moment vector in the external mag. field. one is  $M_I = +\frac{1}{2}$  and other is  $M_I = -\frac{1}{2}$ . In the absence of external mag. field the two mag. moment vectors are degenerate, but as soon as external mag. field is imposed, the degeneracy is lost and it splits into two parts  $M_I = \pm\frac{1}{2}$  in lower energy and as it is produced by the spinning of nucleus in the direction of applied mag. field, the other is  $M_I = -\frac{1}{2}$ , which represents the higher energy level produced by the spinning of the nucleus in the opposite direction of applied mag. field.



If  $\mu$  = mag. moment of the nucleus, and  
 $H_0$  = strength of the applied mag. field.  
 The energy gap b/w the levels is given by

$$\Delta E = -\mu H_0 - (+\mu H_0) = |-2\mu H_0| = 2\mu H_0 \quad \text{--- (1)}$$

As,  $\mu$  for a particular nucleus is constant, the value of  $\Delta E$  depends on the strength of the applied mag. field strength  $H_0$ . We also know that  $\Delta E = h\nu$

From eq. (1)  $h\nu = 2\mu H_0$   
 $\therefore \nu = \frac{2\mu H_0}{h}$  --- (2)

This shows  $\nu$  is proportional to  $H_0$  ( $\nu \propto H_0$ )

Now, if an emt. radiation of radio frequency range is passed through the sample of nuclei populated in lower state may absorb the energy whose  $\nu$  is matching to eq. (2) and it is excited in the higher ~~low~~ state. This is called Flipping of spin or flipping of protons. When,  $\nu = \frac{2\mu H_0}{h}$  the nucleus is said to be in resonance with applied emt radiations.

Since  $\Delta E$  is not very large compared to 'KT' i.e., thermal energy, the thermal agitation reduces the excess nuclei in lower energy state and thus the two states become equally populated at normal temperature. A/c to Boltzmann's distribution law

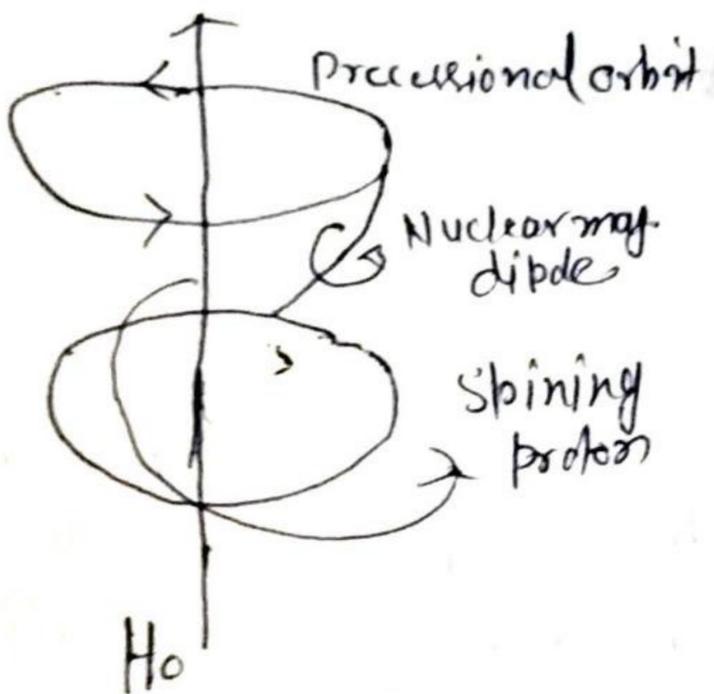
$$\frac{N_e}{N_g} = 1.000006 = e^{-\Delta E/KT}$$

So, when the state of resonance is reached the flipping of the nuclei from  $M_I = +\frac{1}{2}$  to  $M_I = -\frac{1}{2}$  is equally probable and possible for flipping of nuclei from  $M_I = -\frac{1}{2}$  to  $M_I = +\frac{1}{2}$  and subsequently no spectrum is expected.

Q. How does interaction b/w nuclear mag. moment and external mag. field occur?

Ans.:

A nucleus having nuclear spin  $\frac{1}{2}$  (or more) may be assumed as a bar magnet which spins on its own axis. When it is placed in an external mag. field, the interaction b/w the moment of the nuclear bar magnet and the external mag. field produces torque. This torque causes the nuclear mag. moment vector to precess about the applied mag. field vector ( $H_0$ ). The phenomenon is called Larmor precession.



Proton precessing in mag. field  $H_0$

The angular frequency of this Larmor precession is called Larmor frequency,  $\omega$  rad. sec. It depends on the strength of ~~mag~~ field the applied mag. field  $H_0$ ,

$$\omega = \gamma \cdot H_0$$

where  $\gamma$  is magnetic gyric ratio or Gyromagnetic ratio (or constant)

$$\gamma = \frac{2\pi\mu}{hI}$$

Here  $\mu$  = mag. moment of the spinning bar magnet

$I$  = spin Q. No. of the bar magnet nucleus

$h$  = Planck's const

So, at first to the nuclei of spin  $I = \frac{1}{2}$  (or more), at first a strong mag. field is applied externally to cause the nuclei to precess about it.

Now emt. radiation is applied which is comparable to  $\Delta E$  and hence radiofrequency range. When the frequency of applied emt. radiation is just equal to the Larmor frequency of the nucleus, it comes into resonance with the applied emt. radiation.

Under this condition the nuclei may absorb energy from the radiation and flip into the higher energy state or a nucleus in higher energy state may emit energy and flip back to the lower state as both the states are almost equally populated at room temperature. So, a mag. field of stronger strength is applied so that a little more nuclei may be in ground level, and hence under the condition of resonance, the flipping of protons occur from the lower to higher state producing signal in detectors i.e., n.m.r. spectrum.

[ we know that A/c to fundamental 19112 equation which correlates emit. frequencies with mag. field, we say that

$$\omega = 2\pi\nu = \gamma/H_0 \quad ]$$

we know that  $\omega = 2\pi\nu$  and  $\Delta E = h\nu$

So,  $\Delta E \propto \omega$  (frequency of Larmor precession)

So, resonance state is reached only when  $\omega = 2\pi\nu$

where  $\nu$  = radio frequency (r.f.)

This condition may be reached in two different ways :-

(i) Applied mag. field is kept const. and thus Larmor frequency of nuclei is fixed and now the applied r.f. is scanned until the matching occurs i.e.,  $\omega = 2\pi\nu$

(ii) In other way, the radiofrequency is fixed and the strength of applied mag. field i.e.,  $H_0$  is varied until the matching occur i.e.,  $\omega = 2\pi\nu = 2\pi \left( \frac{2\mu H_0}{h} \right)$

The most n.m.r. instrument follows the 2<sup>nd</sup> process. In both the process, when matching occur the flipping of protons start producing signal i.e., n.m.r. spectrum.

However, after initial flipping again both the states are equally populated. Under this state the sample is called saturated.

However, the state of saturation is avoided in the instruments by proper mechanism. In the instruments, that arrangement is made in such a way that, the excited nuclei may return to the ground state no by emitting energy directly, rather by some other mechanism. This happens by two mechanisms

- (i) spin-spin relaxation and
- (ii) spin-lattice relaxation

~~Precessi~~ Precessional Frequency ( $\omega$ )  $\rightarrow$   
It may be defined as the number of revolutions per second made by the mag. moment vector of the nucleus around the external field  $H_0$ .

or  
It may be defined as the precessional frequency of the spinning bar magnet (nucleus) equal to the frequency of ext. radiation in mega cycle per second necessary to induce a transition from one spin state to another.